



Non-trivial discrimination between the Scotogenic model and the Inert Doublet model

(Looking for the difference between
Tweedledum and Tweedledee)



October 18 - IFIC

Who am I?



Antofagasta



Santiago



Concepción

The cover of a thesis. At the top left is the logo of the University of Concepción. The title "Propagation of quantum correlations through a turbid media" is written in a bold, serif font on a pink rectangular background. Below the title, the author's name "Ivana Matorana Avila" is listed, followed by "Advisor: Dr. Carlos Saavedra" and "Co-advisor: Dr. Esteban Sepúlveda" in a smaller font.

The Cast



Sebastian Urrutia



Ivania Maturana

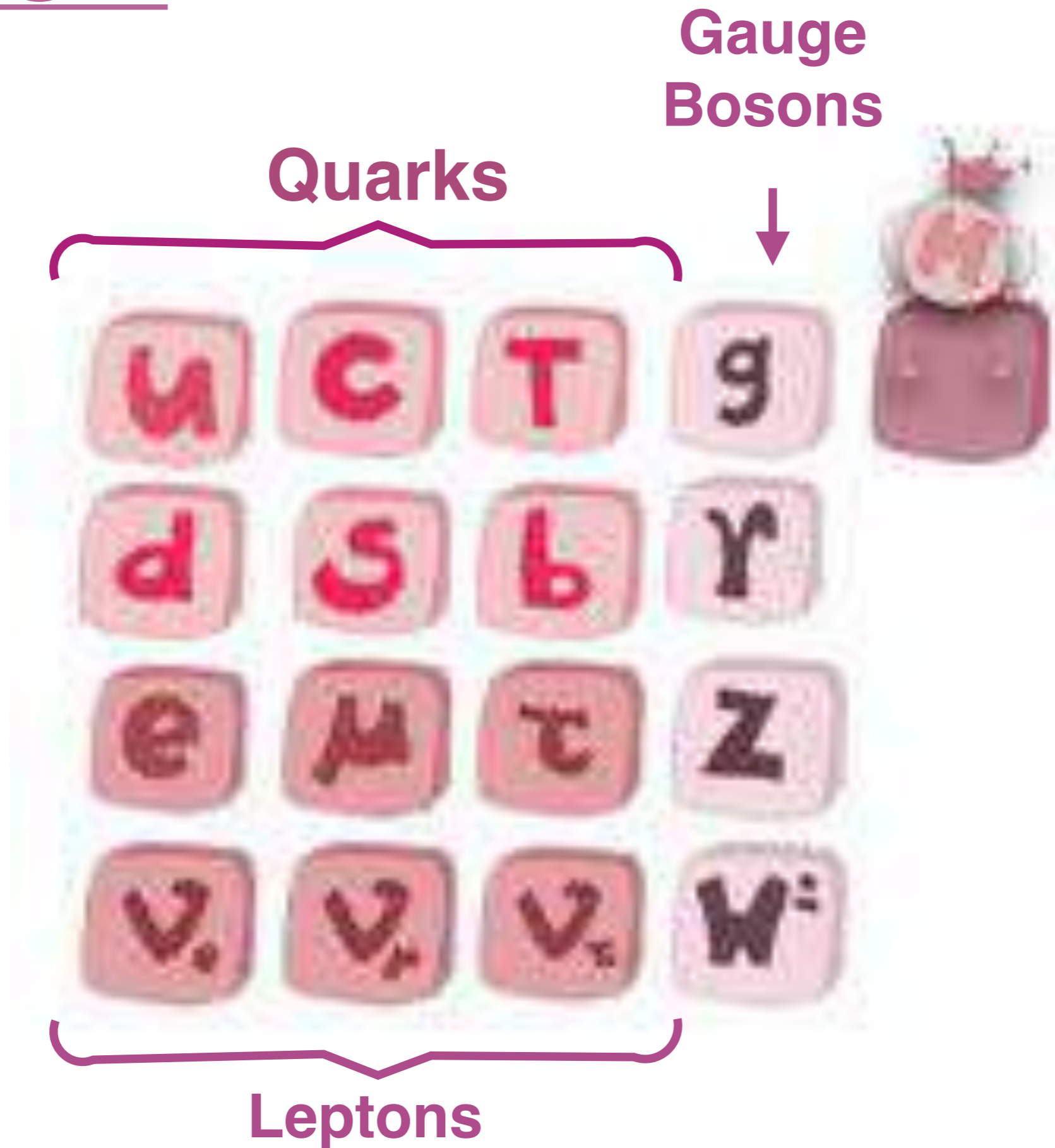


Professor MAD



Nicolás Rojas

The SM



The SM problems



- Is there a theory that describes all the fundamental forces?
- Why are there 3 generations of quarks and leptons with different scales of mass?
- Why Higgs mass is so small?
- What is dark matter?
- How the neutrinos acquire mass?



This is not a work about finding a solution to the SM questions



The SM problems



- Does exist a theory that describes all the fundamental forces?
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- How the neutrinos acquire mass?



We want to find a physical observable that help us to discriminate between both models.



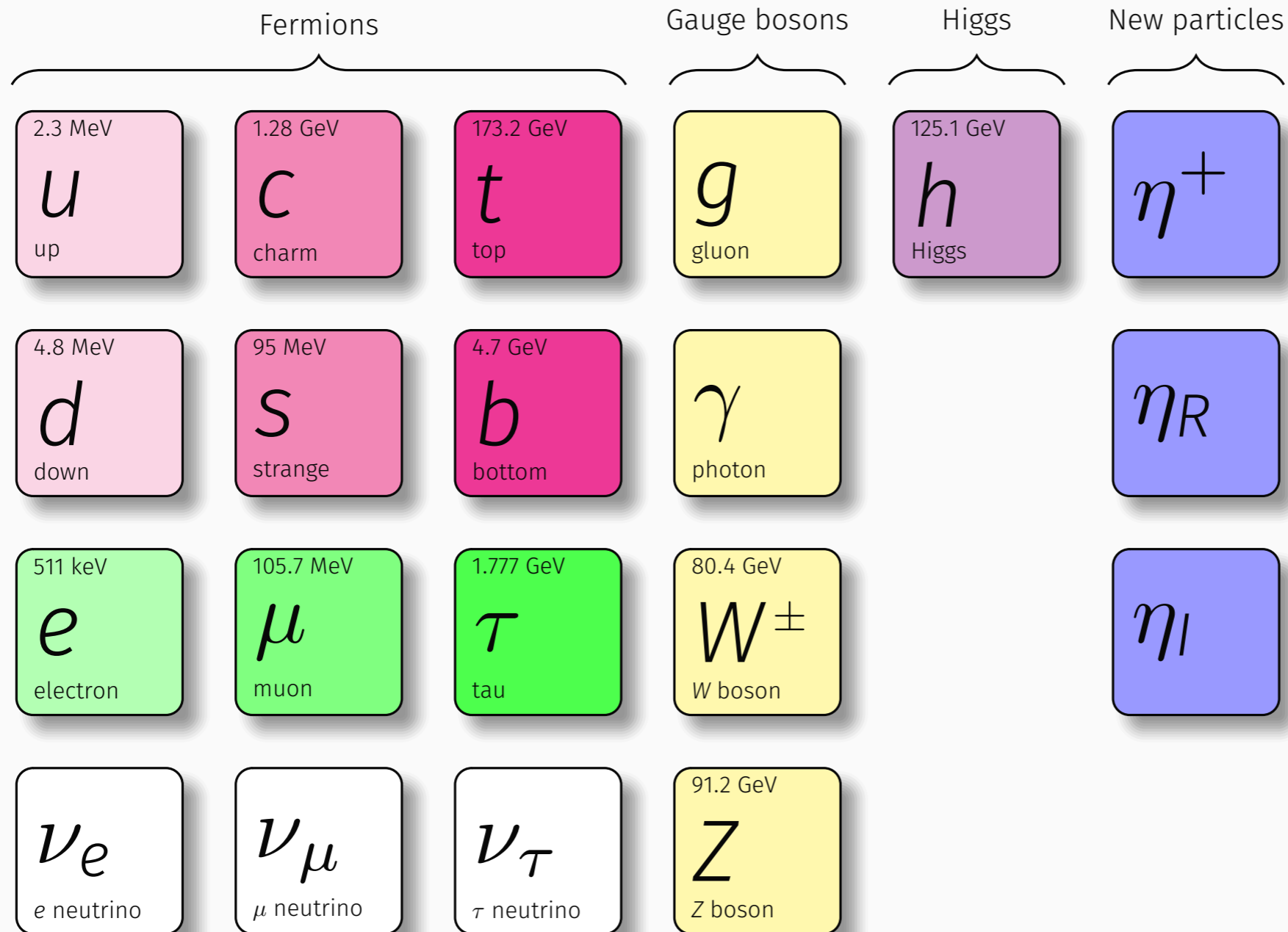
Inert doublet model

VS

Scotogenic Model

Why?

The inert doublet model



We can extend the SM adding two doublets charged under SU(2). The name of this model is THDM (1). The Inert doublet model adds a Z_2 symmetry or a stabilizing symmetry. The new symmetry acts in the new particles of the model as follows

$$\begin{aligned} \eta &\rightarrow -\eta \\ SM &\rightarrow SM \end{aligned} \tag{1}$$

The scalar potential is

$$\begin{aligned} V = & m_1^2 \phi^\dagger \phi + m_2^2 \eta^\dagger \eta + \lambda_1 (\phi^\dagger \phi)^2 + \lambda_2 (\eta^\dagger \eta)^2 + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\ & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) + \frac{\lambda_5}{2} [(\phi^\dagger \eta)^2 + (\eta^\dagger \phi)^2], \end{aligned} \tag{2}$$

where m_1^2, m_2^2 and λ_{1-5} are reals.

Expanding the fields around the vacuum, we will get

$$\begin{aligned} \phi &= \begin{pmatrix} 0 \\ (v+h)/\sqrt{2} \end{pmatrix}, \\ \eta &= \begin{pmatrix} \eta^+ \\ (\eta_R + i\eta_I)/\sqrt{2} \end{pmatrix}. \end{aligned} \tag{3}$$

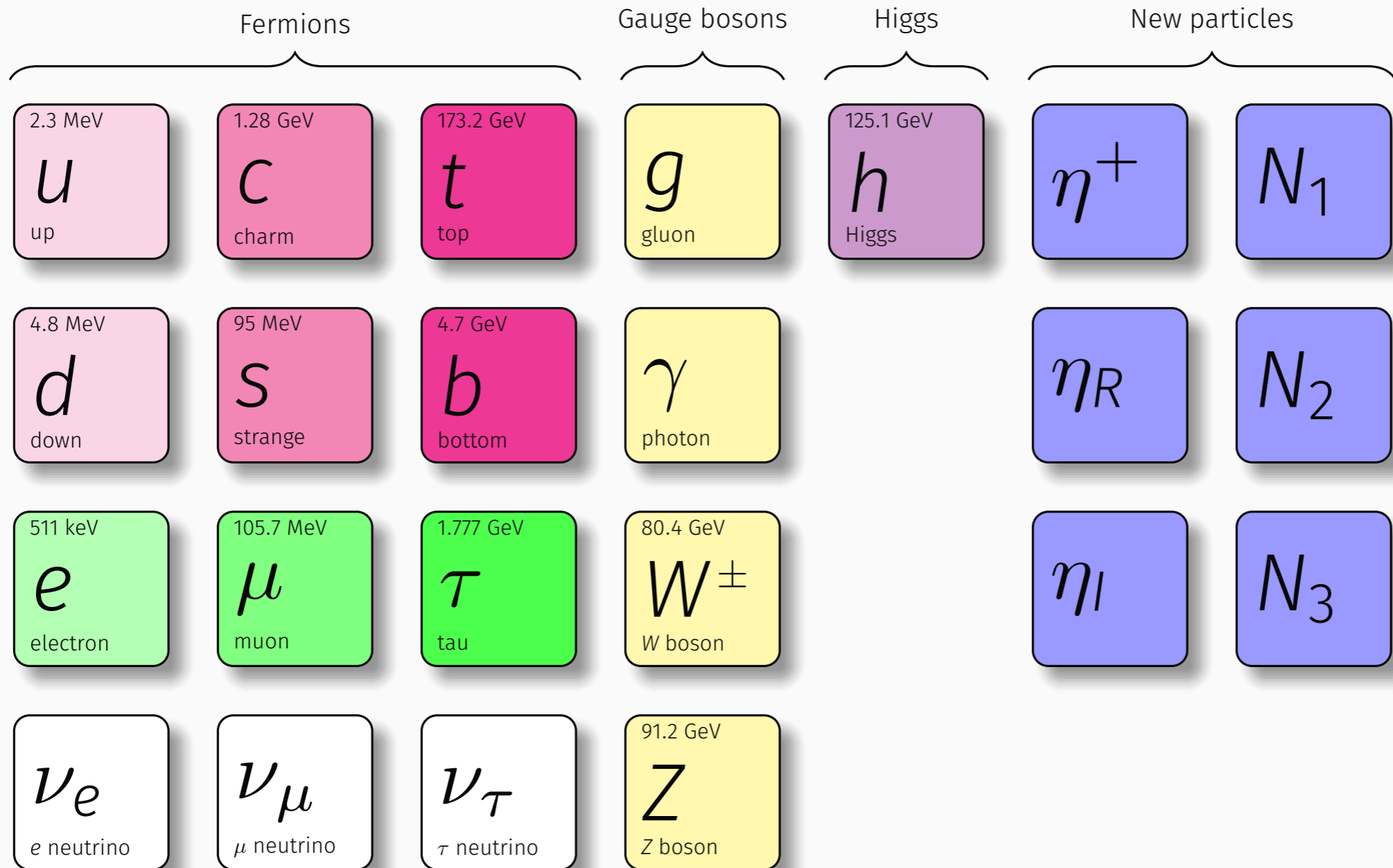
The mass terms of the physical states are

$$\begin{aligned}m_{\phi}^2 &= 2\lambda_1 v^2, \\m_{\pm}^2 &= m_2^2 + \frac{\lambda_3}{2} v^2, \\m_R^2 &= m_2^2 + \frac{\lambda_3}{2} v^2 + \left(\frac{\lambda_4 + \lambda_5}{2}\right) v^2, \\m_I^2 &= m_2^2 + \frac{\lambda_3}{2} v^2 + \left(\frac{\lambda_4 - \lambda_5}{2}\right) v^2.\end{aligned}\tag{4}$$

In addition, the parameters of our model have the next constrains

$$\begin{aligned}\lambda_{1,2} &> 0 \\ \lambda_3 &> -\sqrt{\lambda_1 \lambda_2} \\ \lambda_3 + \lambda_4 - |\lambda_5| &> -\sqrt{\lambda_1 \lambda_2} \\ \lambda_4 + \lambda_5 &< 0 \\ \lambda_5 &< 0\end{aligned}\tag{5}$$

The scotogenic model



The Scotogenic model (1) introduces a new Z_2 symmetry, the same that was introduced in the Inert Doublet model and will act over the fields like

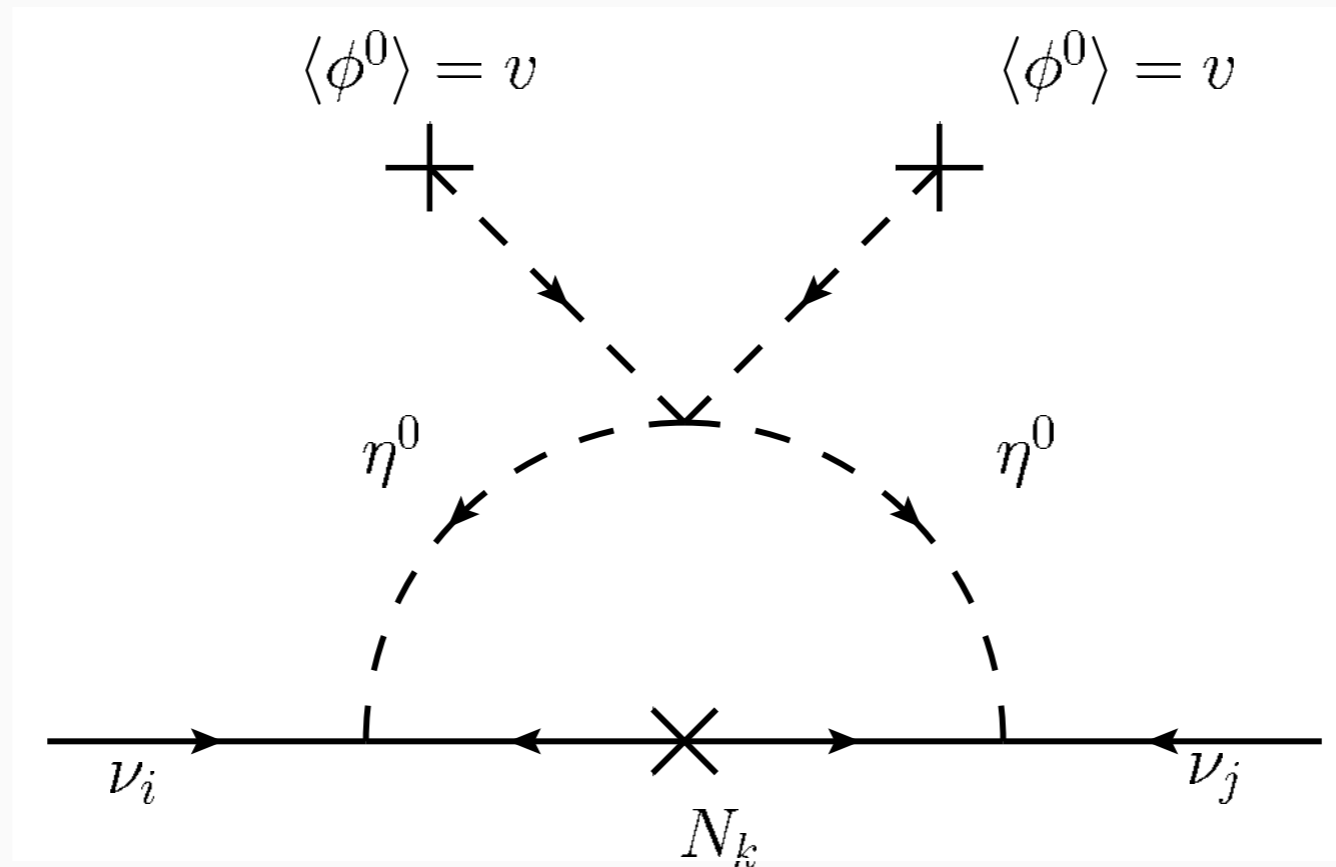
$$\begin{aligned}
 N_{1,2,3} &\rightarrow -N_{1,2,3} \\
 \eta &\rightarrow -\eta \\
 SM &\rightarrow SM
 \end{aligned}
 \tag{6}$$

The model will generate neutrino masses at one loop. Comparing the model with the SM, the new fields/terms that will appear in the Lagrangian are

1. $\eta \approx (2, 1/2)$.
2. $N_i \approx (1, 0)$, that will have a Majorana mass term given by $-\frac{1}{2} \tilde{N}^i M_{ij} N^{jc}$, with M_{ij} a matrix.
3. The Yukawa coupling $L_{Yukawa} \supset -h_{ij} \bar{N}_R^i \tilde{\eta}^\dagger l_L^j + h.c.$ (with $\tilde{\eta} = i\sigma_2 \eta^*$).

The neutrino masses will be given by the new term

$$M_{\nu,ij} = \frac{h_{ik}h_{jk}}{16\pi^2} N_k \left[\frac{m_R^2}{m_R^2 - N_k^2} \ln \frac{m_R^2}{N_k^2} - \frac{m_l^2}{m_l^2 - N_k^2} \ln \frac{m_l^2}{N_k^2} \right]. \quad (7)$$



Radiative seesaw mechanism



We found the RGEs (1) for the Scotogenic model, where the most important for us is

$$Dm_2^2 = 6\lambda_2 m_2^2 + 2(2\lambda_3 + \lambda_4)m_1^2 + m_2^2 \left[2T_\nu - \frac{3}{2}(g_1^2 + 3g_2^2) \right] - 4\text{Tr}(M_N M_N^* h h^\dagger), \quad (8)$$

where $T_\nu = \text{Tr}(h^\dagger h)$.

The last equation shows that the parameter m_2^2 could acquire negative values. If that happens the Z_2 symmetry will be broken (2). Taking in account some considerations we realized and study of the Z_2 breaking considering different mass values for N and η :

1.- The mass of the light neutrinos, the vev and the mixing angles will not change so much in the different scales of energy, where the fermionic fields N are introduced.

2.- Firstly we considered $[M_1, M_2, M_3] = [900, 1500, 5000]$ GeV.

(1) [A. Merle, M. Platscher. 10.1103/PhysRevD.92.095002](https://arxiv.org/abs/10.1103/PhysRevD.92.095002)

(2) [A. Merle, M. Platscher. 10.1007/JHEP11\(2015\)148](https://arxiv.org/abs/10.1007/JHEP11(2015)148)



The steps were:

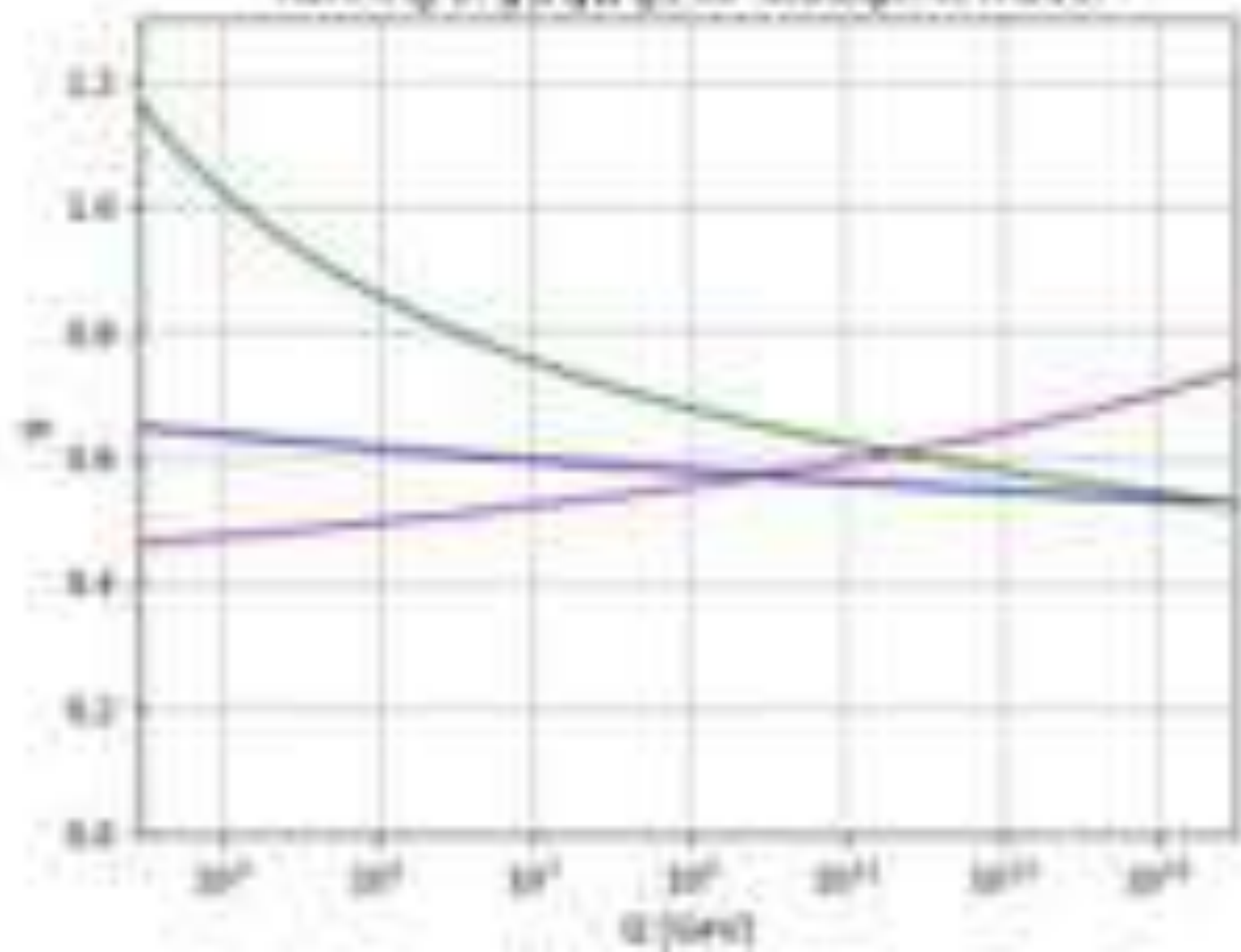
- We calculated the values of the coupling h with a parametrization (1) at the scale of 5000 GeV.

$$h = \sqrt{\Lambda}^{-1} R \sqrt{\hat{m}_\nu} U^\dagger \quad (9)$$

$$\Lambda_i = \frac{M_{N_i}}{(4\pi)^2} \left[\frac{m_R^2}{m_R^2 - M_{N_i}^2} \log \left(\frac{m_R^2}{M_{N_i}^2} \right) - \frac{m_l^2}{m_l^2 - M_{N_i}^2} \log \left(\frac{m_l^2}{M_{N_i}^2} \right) \right]$$

- The parameters $g_1, g_2, g_3, Y_e, Y_u, Y_d, \lambda_1$ and m_1^2 were evolved from 90 GeV scale up to 5000 GeV scale where we save their values.
- Having all the RGEs of the model, we evolved it from 5000 GeV to GUT scale and we go back from 5000 GeV to 90 GeV.

Running of g_1, g_2, g_3 for isotropic model



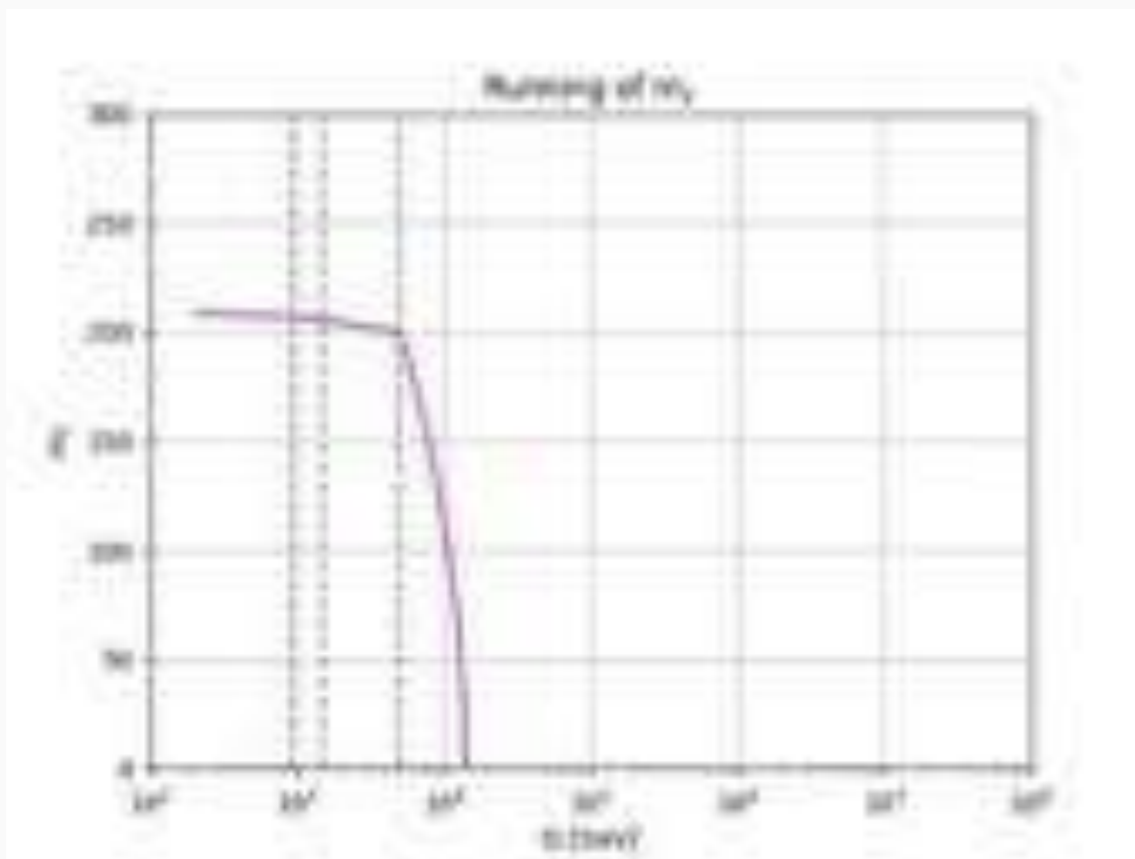


Figure: Plot for $m_2 = 200$ GeV

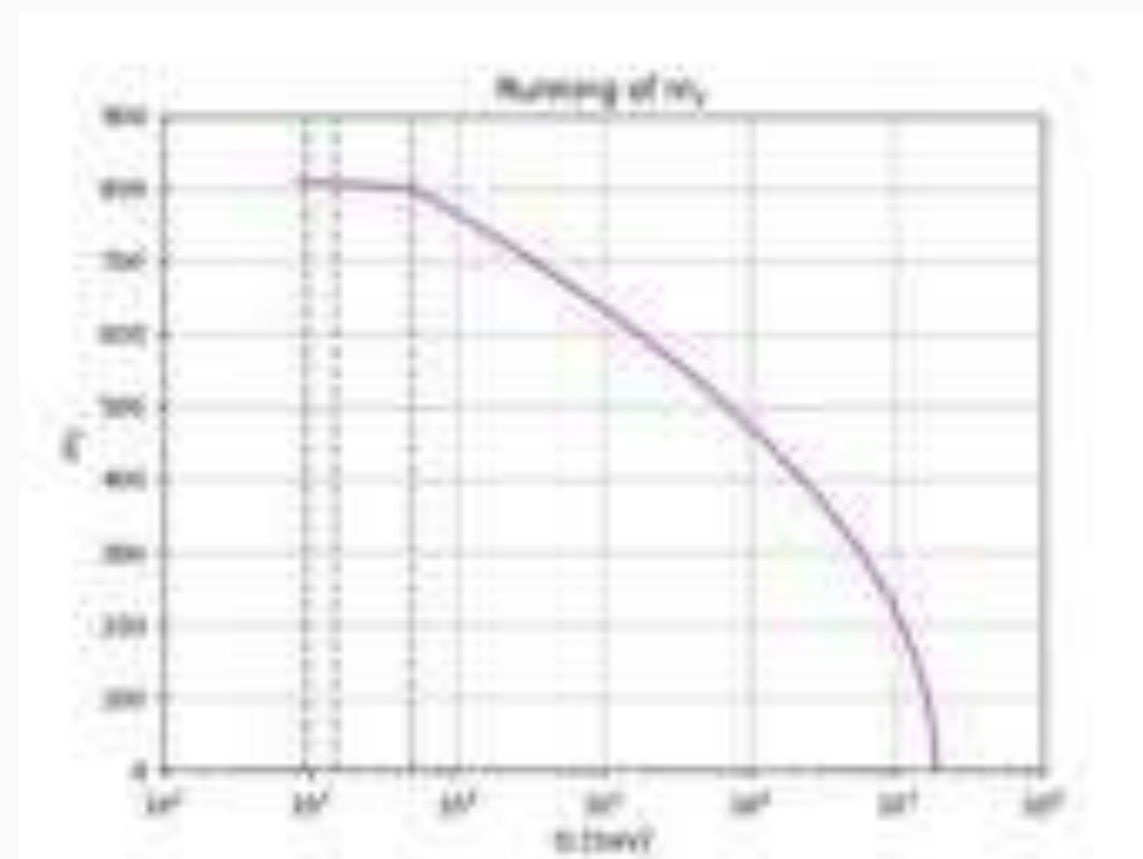


Figure: Plot for $m_2 = 800$ GeV

Changing the mass values for the fermions N and fixing the mass $m_2 = 800$ GeV, we have

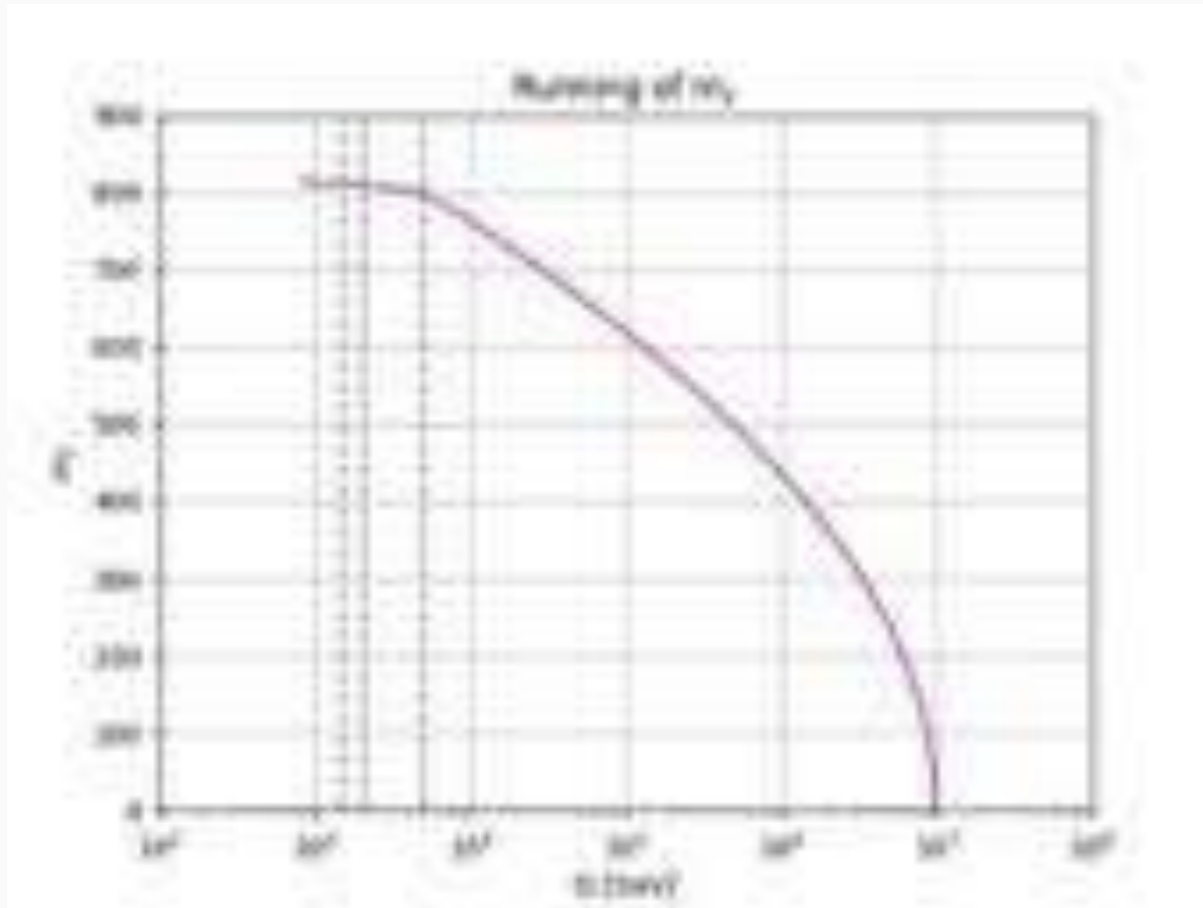


Figure: Plot when $[M_1, M_2, M_3] = [1500, 2000, 5000]$ [GeV]

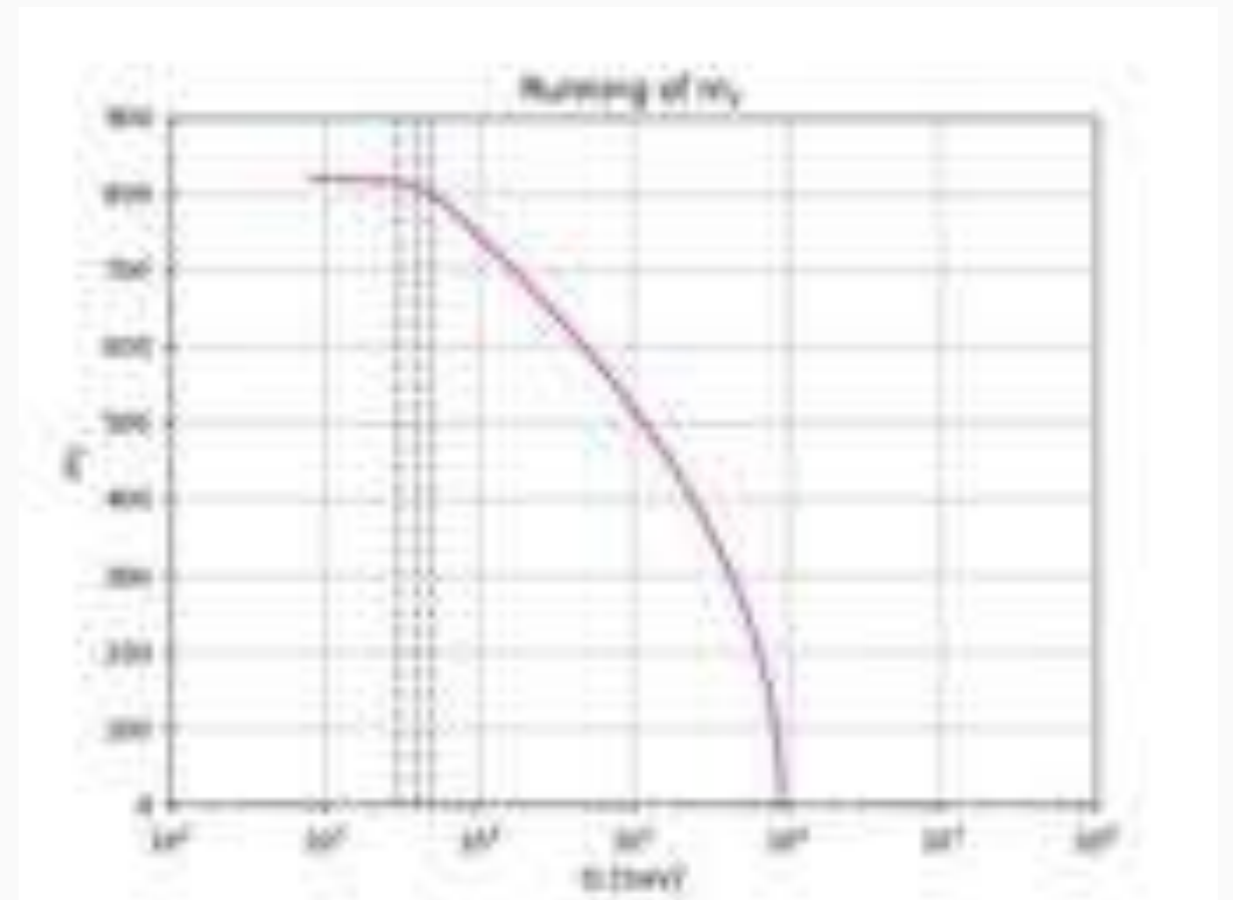
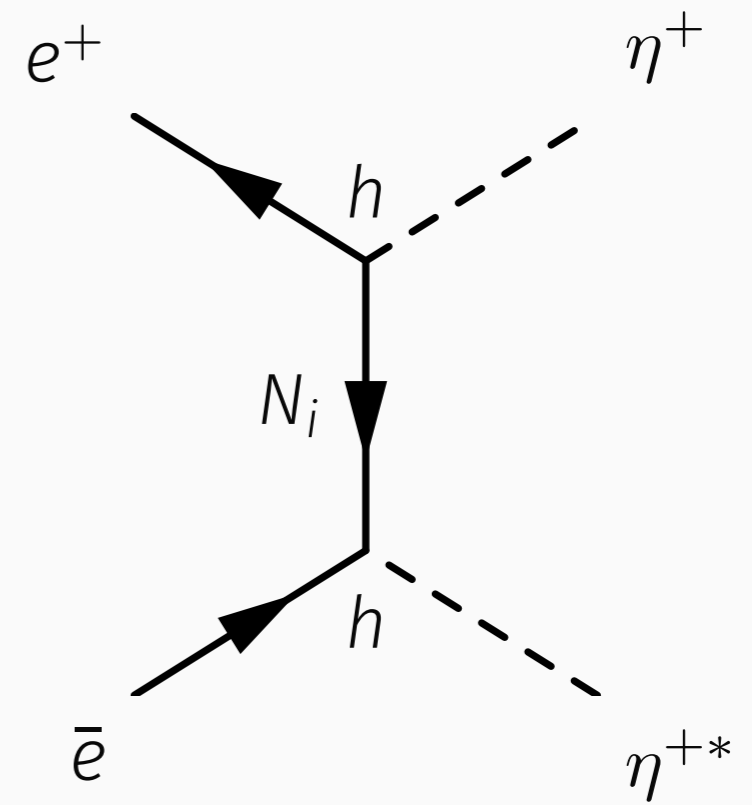
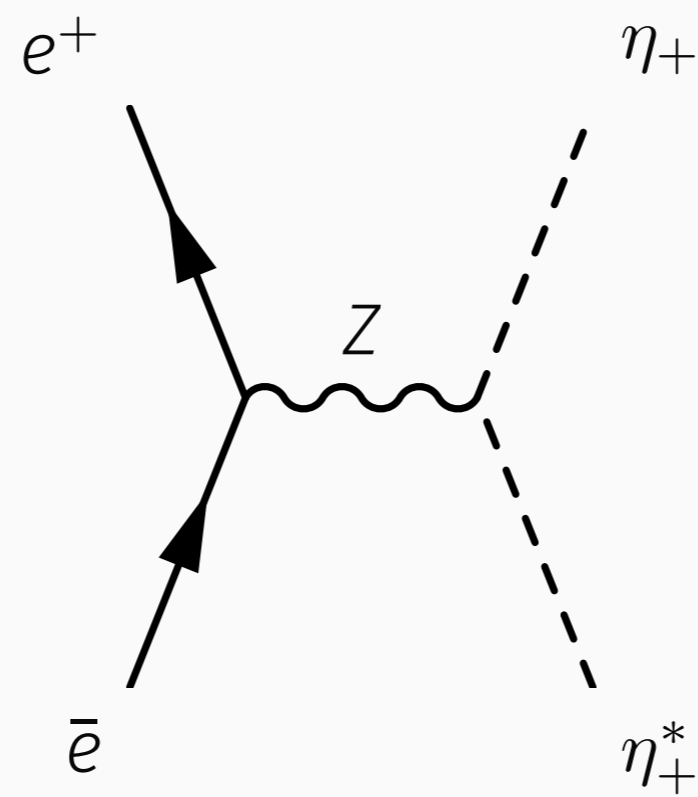
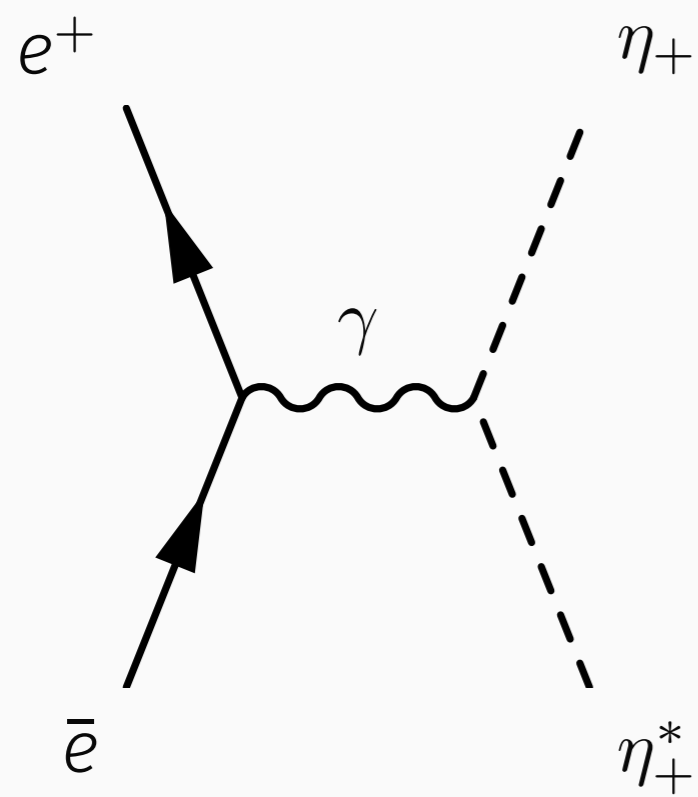


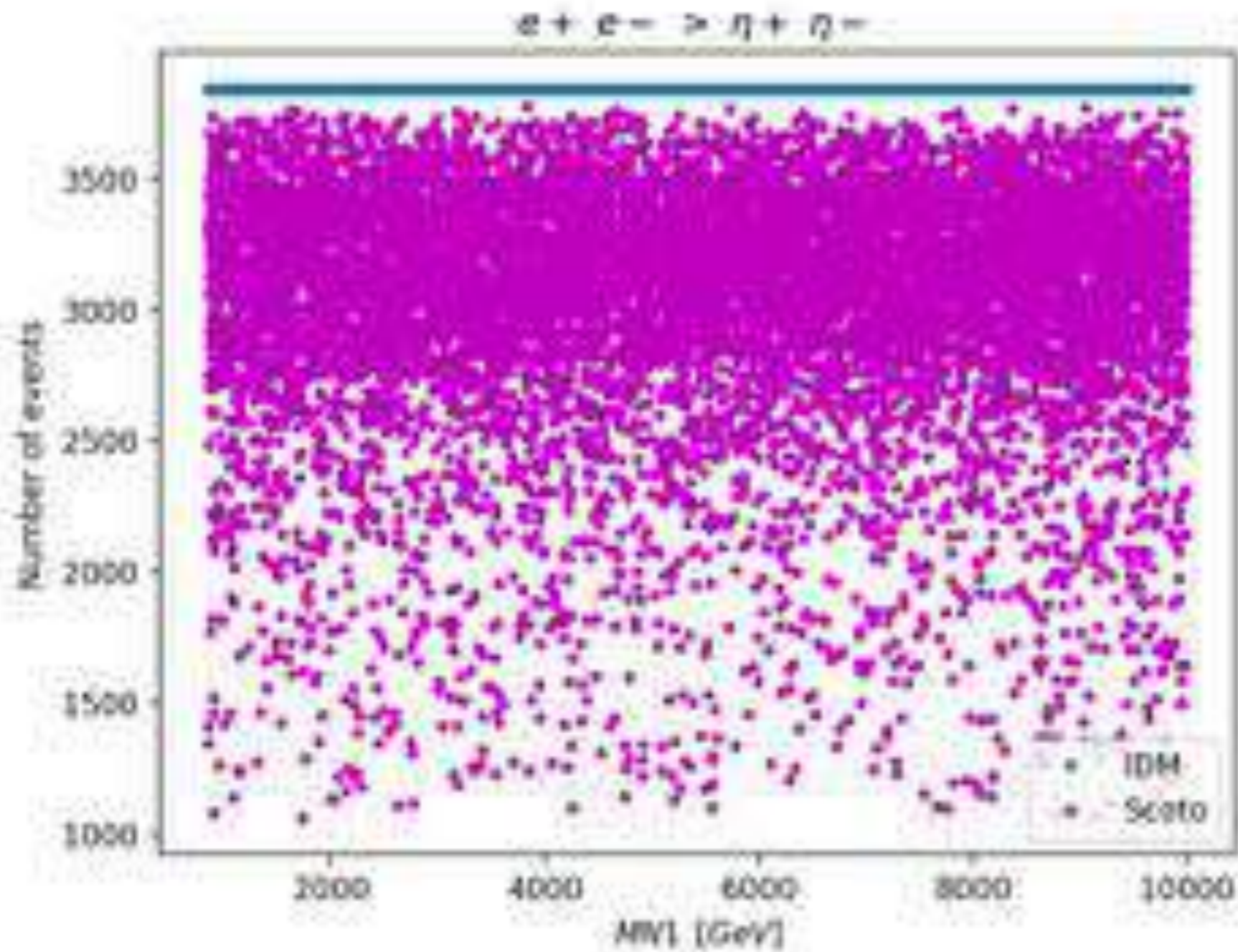
Figure: Plot when $[M_1, M_2, M_3] = [3000, 4000, 5000]$ [GeV]

Finding the observable



$$\sigma_{e^+e^- \Rightarrow \eta^+\eta^-}$$





$$MN_1, MN_2 = [850, 10000] \text{ GeV}$$

$$MN_3 = 10000 \text{ GeV.}$$

$$m_2 = 800 \text{ GeV.}$$

$$m_R = 799,99 \text{ GeV.}$$

$$m_l = 800,00 \text{ GeV.}$$

$$m^+ = 801,88 \text{ GeV.}$$

$$m_{\nu_1}^2 = [0, 0, 001] \text{ eV.}$$

$$\lambda_1 = 0,26$$

$$\lambda_2 = [0, 0,5].$$

$$\lambda_3 = 0,1.$$

$$\lambda_4 = -0,1.$$

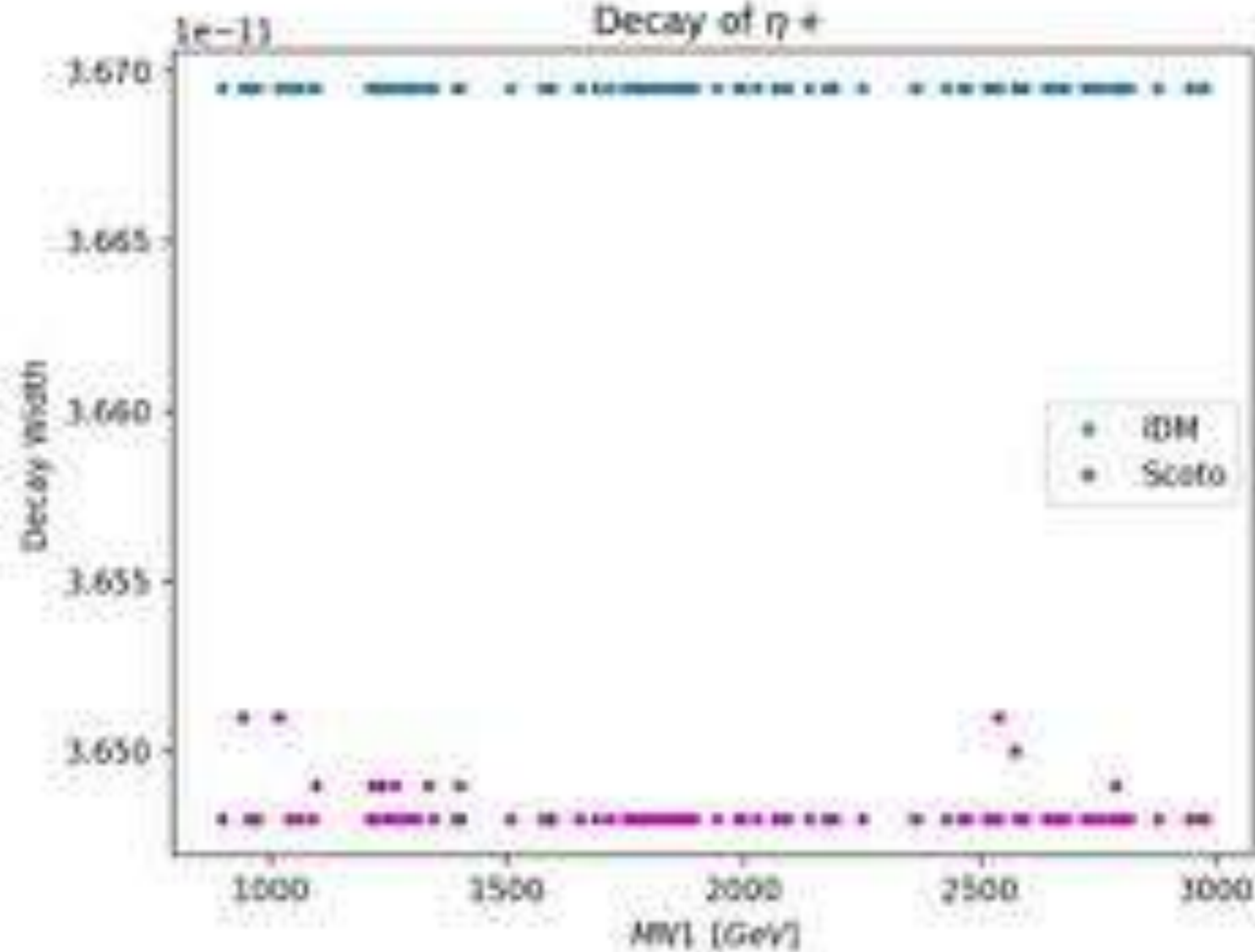
$$\lambda_5 = 10^{-9}.$$

Thanks!

$\sqrt{m^2c^2 + p^2}$
 $SU(3) \times SU(2) \times U(1)$
 $\vec{\nabla} \cdot \vec{D} = \rho$
 $(i\hbar - m)\psi = 0$
 $dU = TdS - PdV$
 $\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$
 \vec{E}
 $\vec{\nabla} \cdot \vec{B} = 0$
 $\sigma_x \sigma_p \geq \frac{\hbar}{2}$
 ϵ_0
 $y(t) = y_0 + v_0 t + \frac{1}{2} a t^2$
 $S = k_B \ln \Omega$
 $\epsilon/\omega \delta b = \delta$
 $\oint \vec{E} \cdot d\vec{S} = qV/\epsilon_0$
 $\vec{F} = q\vec{E}$
 $C = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$
 \hbar
 $\mathcal{L} = i\hbar c \bar{\psi} \not{\partial} \psi - m c^2 \bar{\psi} \psi$
 $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$
 $\vec{r} = \vec{r}/t$

T-U



Decay of $\eta \rightarrow \pi^0 \pi^0$ 

Número Leptónico (L)

$$\left. \begin{aligned} L_{SM} &= 0 \\ L_{\phi} &= 0 \\ L_{\eta} &= 0 \\ L_N &= 1 \end{aligned} \right\} -\frac{1}{2} \bar{N}_R^i M_{ij} N_R^{jc} + h.c.$$

$$\left. \begin{aligned} L_{SM} &= 0 \\ L_{\phi} &= 0 \\ L_{\eta} &= 0 \\ L_N &= 0 \end{aligned} \right\} L_Y \supset -h_{ij} \bar{N}_R^i \tilde{\eta}_\dagger^j l_L^j + h.c$$

$$\left. \begin{aligned} L_{SM} &= 0 \\ L_{\phi} &= 0 \\ L_{\eta} &= 1 \\ L_N &= 0 \end{aligned} \right\} \frac{\lambda_5}{2} [(\eta^\dagger \phi)^2 + h.c.]$$

The SM “problems”



- Does exist a theory that describes all the fundamental forces?

Gravity?

